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NMR OF ²³Na IN NATROLITE

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The temperature dependences of NMR and MAS NMR spectra of ²³Na nuclei in Al-natrolite (Na₂Al₂Si₃O₁₀·2H₂O) and Ga-natrolite (Na₂Ga₂Si₃O₁₀·2H₂O) have been studied. It has been shown that the diffusion of the sodium ions at T < 400 K is absent in Al- and Ga- natrolites. The temperature dependences of the spin-lattice relaxation times T_1 in Al- and Ga- natrolite have been studied. The influence of the water molecular mobility in the nanochannels of natrolites on the spin-lattice relaxation times of ²³Na has been discussed.

Keywords: NMR, magnetic relaxation, zeolite, natrolite, water mobility.

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INTRODUCTION

The Al-natrolite (Na₂Al₂Si₃O₁₀ ·2H₂O) and Ga-natrolite (Na₂Ga₂Si₃O₁₀ ·2H₂O) are the typical channel-type compound with porous structure [1]. The natrolite framework consist of the chains AlO₄ and SiO₄ tetrahedra linked together via common oxygen atoms. The natrolite structure contains channels running both perpendicular and parallel to the c-axis. The water molecules and ions Na+ are located in the small nanochannels of natrolite in the form of zig-zag chains [1, 2] Each sodium ion is coordinated by two framework's oxygen atoms and by two water molecules. The sodium cations in the channels of natrolite, as well as water molecules, can possess the certain mobility. However unlike molecules of water they cannot be removed from a crystal if only to not resort by an ionic exchange [2].

The dynamics of water molecule in Al-natrolite have been studied by NMR in [2-4]. From temperature measurements of the spin-relaxation times ($T_1 \ \mu \ T_{1\rho}$) of ¹H nuclei in Alnatrolite it was concluded that in an temperature interval 330 K ÷ 450 K takes place the reorientation of water molecules around of their pseudo-axes of second order symmetry (180° flip motion), and in an interval 450 K ÷ 540 K it has been assumed availability reorientation of water molecules around of the one hydrogen bonds [3]. In the subsequent it has been established, that this second type of water molecules mobility is connected with diffusion of water molecules in the channels of natrolite and the 180° flip motion take place simultaneously with diffusion along the c-axis [2, 4].

The powerful method for the study of zeolites is the NMR of quadrupole nuclei. The NMR of the quadrupole nuclei ²³Na in a single crystal of Al-natrolite have been investigated at room temperature in [5, 6]. The obtained quadrupole coupling constant (QCC) and asymmetry parameter of the electric field gradient (EFG) tensor were determined to be: eqQ/h = 1759 kHz; $\eta = 0.6427$. From results represented in [4] it

follows that the QCC of ²³Na nuclei do not depend on the temperature at T < 500 K. However at T > 500 K the QCC are decreased and the observed decreasing of the QCC of ²³Na nuclei is connected with diffusion of water molecules in the natrolite channels. From the observed temperature independent of NMR spectra of ²³Na in Al-natrolite single crystal it was concluded that the translation diffusion of Na ions is absent in the natrolite cannels [4].

In the present paper we represent the results of the study of Al- and Ga-natrolites by the ²³Na NMR. We investigate the temperature dependences of the ²³Na NMR spectra in static and rotated (MAS) samples with and without ¹H decoupling as well as the temperature dependences of the spin-lattice relaxation times of the ²³Na nuclei in Al- and Ga-natrolites.

1. EXPERIMENTAL PART

The polycrystalline samples of natural Al-natrolite from Khibiny deposit (Kola Peninsula, Russia) were used in this study. The gallium form of natrolite was hydrothermaly synthesized as described in [7]. The ²³Na NMR spectra were measured at $v_0 = 105.842$ MHz frequency in 9.4 T magnetic field using a Bruker Avance-400 NMR spectrometer. The ²³Na MAS NMR spectra were measured using the 4 mm diameter zirconia (ZrO₂) rotor cells equipped with powdered sample rotated under magic angle with 10 kHz frequency. Classical direct acquisitions by single pulse excitations (free induction decay - FID) were used. The ²³Na has I = 3/2 nuclear spin, and for selective excitation of the central ($m_I = +1/2 \leftrightarrow -1/2$) transition the optimal pulse duration equals to the duration of a non-selective $\pi/2$ pulse divided by I + 1/2 = 2 [8]. In our experiments the radiofrequency (RF) puls $\pi/4 = 1.0$ µs was used. The NMR spectra were obtained by Fourier transformation of FID signals. The spin-lattice relaxation time T_1 for ²³Na nuclei was measured by saturation-recovery method.

The Dmfit program [9] was used to simulate the ²³Na spectra to extract the isotropic chemical shifts (δ_{iso}), quadrupolar coupling constants and the asymmetry parameters (η_Q). The Dmfit model of MAS NMR spectrum includes an apodisation of the theoretical lineshape by Lorentzian or Gaussian curves with the broadening parameters Δv , Δv_L or Δv_G , that indicate a distribution of slightly distinct environment of nuclei.

2. OBTAINED RESULTS AND DISCUSSION

The sodium neighbours in Al- and Ga-natrolites have a configuration of a distorted tetrahedron. In the tetrahedron corners there are two oxygen atoms belonging to a framework and two oxygen atoms of water molecules at an average distances of 2.37 Å. Furthermore there are two oxygen atoms of a framework at an average distances of 2.5 Å, four protons at an average distances of 2.8 Å and atoms of silicon and aluminum at an average distances of 3.0 Å.

The experimental ²³Na MAS NMR spectrum of Al-natrolite is shown in Fig. 1, *a*. The rotation of sample at magic angle (MAS NMR) leads to fully average of dipolar interaction of magnetic moments of ²³Na with magnetic moments of other nuclei in natrolite ¹H, Si, Al. In this case the shape of NMR spectra is determined only by the second-order quadrupolar shift of the central transition [8]. Using program DMFit [9] we calculated the shape of MAS NMR spectrum of ²³Na in polycrystalline natrolite. The result of theoretical calculations are shown in Fig. 1, *b*. The obtained theoretical values of the quadrupolar frequency v_Q and the asymmetry parameter η well coincide with experimental values obtained in [5].

In Fig. 2 is presented the ²³Na MAS NMR spectra obtained in Al-natrolite at T = 300 K and T = 380 K. These spectra were obtained using the method of ¹H decoupling. NMR ¹H decoupling is a special method used in NMR which allow to eliminate fully the effect of magnetic dipolar coupling between resonance nuclei (²³Na in our case) and ¹H nuclei. From the comparison of the ²³Na NMR spectra, shown in Fig. 2, it appears that the ²³Na NMR spectra have the same shape at T = 300 K and T = 380 K. The NMR shape of ²³Na is determined by magnetic dipolar interactions with other magnetic nuclei and by electric quadrupolar interaction with the electric field gradient (EFG) on the ²³Na sites. The interaction with magnetic moment of ¹H nuclei give the main contribution to the magnetic dipolar interaction of ²³Na nuclei. However the ¹H decoupling, which was used at recording the NMR spectra of ²³Na NMR (Fig. 2), leads to averaging of the dipolar interactions between the magnetic moments of ¹H and ²³Na nuclei and so, from Fig. 2, it follows that in the temperature region T < 380 K the electric field gradient (EFG) at the ²³Na sites does not depend on the temperature.



Fig. 1. ²³Na MAS NMR spectra of Al-natrolite at $\Omega_{rot} = 10$ kHz: (a) experimental spectrum at T = 300 K; (b) theoretical spectrum with parameters $C_Q = 1759,2$ kHz, $\eta = 0,64, \delta_{CSA} = 8,19$ ppm, $\Delta v_{Gauss} = 73,11$ Hz.



Fig. 2. ²³Na NMR spectra with ¹H-decoupling in Al-natrolite at T = 300 K and T = 380 K.

In natrolite there are two structurally nonequivalent ²³Na and well-resolved fine structure of NMR spectra caused by the second-order qudrupolar effects is observed at some orientations of crystal in external magnetic field [6]. The experimental temperature dependencies of the quadrupolar second-order shift of the central NMR lines of the two structurally nonequivalent ²³Na nuclei, for the case when the vector of the magnetic field \mathbf{B}_0 is parallel to [110] direction, are shown in Fig. 3 [4]. This result was obtained using the CW NMR spectrometer at a frequency of the of 11 MHz on Al-natrolite single crystals [4].



Fig. 3. The temperature dependencies of the second-order quadrupolar shifts and the linewidth (δv) of ²³Na NMR spectra for the two structurally nonequivalent ²³Na ions in the natrolite single crystal. Inset: the central part of NMR spectrum of ²³Na in single crystal.

The intensive diffusion of ²³Na cations in natrolite pores should lead to the averaging of the second-order quadrupolar shifts of the two structurally nonequivalent ²³Na nuclei. In reality, such effect is not observed (Fig. 3). So, from this result it follows that the diffusion of the sodium ions is absent in natrolite channels in the temperature region T < 500 K [4].

The EFG at the ²³Na sites in the natrolite structure is determined by the electric charges of the ions of whole lattice and by the electric dipolar moments of the water molecules. According with ¹H NMR data of Al-natrolite the water molecules at T < 380 K rotate about their quasi 2-fold axis [3] and the 180° flip motion take place simultaneously with diffusion along the c-axis [2]. These motions of the water molecules must lead to the averaging of the contributions of the electric dipolar moments of water molecules to EFG at the ²³Na sites. If contribution from the electric dipolar moments of water molecules to the EFG tensor at the ²³Na sites is considerable the averaging of this contribution must be observable in the temperature dependence of ²³Na NMR spectra. From our experimental temperature dependences of ²³Na MAS NMR and NMR spectra represented in Fig. 2 and Fig. 3 it follows that this effect is not observed. From this fact we may conclude that the contributions of the electric dipolar molecules to EFG at the ²³Na sites are very small.

Fig. 4. shows the ²³Na NMR spectra obtained without ¹H decoupling at T = 300 K and T = 380 K. The difference between the NMR spectra represented in Fig. 2 and Fig. 4 are related to the dipolar interactions between the magnetic moments of ¹H and ²³Na nuclei. The shape of these NMR spectra is determined not only by the second-order quadrupolar shift of the central transition but also by the dipolar interaction between magnetic moment of ²³Na and ¹H nuclei. From Fig. 4 it follows that increasing of the sample temperature leads to thermal averaging of dipolar interactions between the magnetic moments of ¹H and ²³Na nuclei. Because from NMR data it follows that the 180° flip motion of water molecules take place simultaneously with diffusion along the c-axis [2] we may conclude that the averaging of dipolar interactions of the ²³Na with magnetic moments of protons is connected with the rotations of water molecules about their quasi 2-fold axis and with diffusion of water molecules across canals in natrolites [4].



Fig. 4. ²³Na NMR spectra without ¹H-decoupling in Al-natrolite at T = 300 K and T = 380 K.

The temperature dependence of the spin-lattice relaxation times T_1 of ²³Na in Ganatrolite is shown in Fig. 5. The similar temperature dependence is observed in Alnatrolite.

In order to identify the main causes of the measured ${}^{23}Na$ relaxation times in natrolite, we consider the theoretical calculations of T_1 on the basis of different dynamical model. The physical mechanisms which could be induce the longitudinal relaxation of the ${}^{23}Na$ nuclei are [10]:

(i) spin-phonon interactions – interactions of the quadrupolar electric moment of the ^{23}Na nuclei with the crystal electric field gradient modulated by lattice vibrations;

(ii) dipolar interaction with paramagnetic impurities (for example with Fe^{3+});

(iii) magnetic dipolar interaction with absorbed ions or molecules and with magnetic moments of ²⁷Al (⁷¹Ga, ⁶⁹Ga), ²⁹Si, ¹H and other ²³Na nuclei of natrolite structure;

(iv) electric quadrupolar interactions with the crystal electric field gradients modulated by motion of charge cations or water molecules [10]. These motions are the "hopping" motions, i.e. the atoms or water molecules spend most of their time in potential well corresponding to equilibrium positions, and only a very small fraction move between these potential wells.

The phonon-based relaxation mechanism could not be the cause of spin-lattice relaxation of the quadrupolar nuclei in natrolites. From obtained estimations it follows that if relaxation were to proceed by this mechanism the values of T_1 for quadrupolar nuclei in zeolites at room temperature should be 4-5 orders of magnitude larger than the experimental values [10]. The spin-lattice via paramagnetic impurities may be significant only at very low temperature [10]. The magnetic dipolar interaction could not be also the cause of the spin-lattice relaxation of the ²³Na nuclei in natrolite. From our estimations we obtain that the dipolar interactions of the ²³Na with magnetic moments of the proton magnetic moment give the minimal value of the spin-relaxation time T_{1min} equals ~ 5 s, which much larger than experimental values 0,04 s (Fig. 5). So only one mechanism (iv) need to be considered.

It is known that the relaxation of the quadrupolar nuclei may be multiexponential [11, 12]. However, for selective saturation of the central transition [8], the relaxation is well described by single exponential [13]. In our case the translational and reorientationa jumps of water molecules modulate only the part of EFG tensor concerned with electric dipolar moments of water molecules. The remaining part of the EFG tensor given by the electric charges of the ions of whole lattice is not changed. So the quadrupolar relaxation concerned with modulation of the part of EFG at the site of the quadrupolar nuclei as a result of the activated translational and reorientational jumps of electric dipoles of water molecules may be described (I = 3/2) by equation [10, 12, 13]

$$T_{1Q}^{-1} = 9 \left(1 + \frac{\eta^2}{3} \right) \Delta C_Q^2 \frac{\tau_c}{1 + \omega_0^2 \tau_c^2} , \qquad (2)$$

where ω_0 is the Larmor frequency of the quadrupole nucleus; η – asymmetry parameter of EFG tensor and



Fig. 5. The temperature dependences of the spin-lattice relaxation time of the 23 Na nuclei in Ga – natrolite.

$$\Delta C_Q = \frac{e^2 \Delta q Q}{\hbar} \ . \tag{3}$$

Here $e\Delta q$ describes the part of EFG at the site of the ²³Na nuclei connected with the electric dipoles of water molecules.

From Eq. (2) it follows that minimal value of T_{1Q} for selective saturation and detection of the central transition is observed at $\omega_0 \tau_c = 1$ and is equal

$$T_{1\min} = \frac{4\pi v_0}{3(3+\eta^2)\Delta C_Q^2} \ . \tag{4}$$

For the ²³Na nuclei in the natrolite $T_{1\min} \cong 0.04$ s. Using $\eta = 0,64$ and $v_0 = 105,542$ MHz we obtain from Eq. (4)

$$\Delta C_o \approx 56,93 \text{ kHz.}$$
(5)

The full constant of quadrupolar interaction $C_Q = 1759,3$ kHz and so from our result it follows that the contributions of the electric dipolar moments of the water molecules to full EFG at the ²³Na sites are 3.2% only.

CONCLUSION

From analysis of the temperature dependencies of NMR and MAS NMR spectra of ²³Na nuclei (with and without ¹H-decoupling) it follows that the diffusion of the sodium ions at T < 400 K absents in Al- and Ga- natrolites. From analysis of NMR spectra of ²³Na nuclei (with and without ¹H-decoupling) it follows that in Al- and Ga-natrolites the 180° flip motion of water molecules take place simultaneously with diffusion of the water along the Schottky defects. The obtained from ²³Na MAS NMR spectrum theoretical values of the quadrupolar frequency v_Q and the asymmetry parameter η well coincide with experimental values obtained early. The spin-lattice relaxation of the ²³Na is governed by the electric quadrupole interaction with the crystal electric field gradients modulated by translational motion of H₂O molecules in the natrolite pores. The dipolar interactions with paramagnetic impurities become significant as a relaxation mechanism of the ²³Na nuclei only at low temperature (< 270 K).

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В Al-натроліті (Na₂Al₂Si₃O₁₀ · 2H₂O) та Ga-натроліте (Na₂Ga₂Si₃O₁₀ · 2H₂O) вивчені температурні залежності спектрів ЯМР і МАЅ ЯМР ядер ²³Na. Було показано, що в Al-натроліті та Ga-натроліті при T < 400 К відсутня дифузія іонів натрію. Отримано температурні залежності часу спін-граткових релаксації T_1 в Al-натроліті і Ga-натроліті. Обговорюється вплив молекулярної рухливості води в наноканалах структури натроліта на спин-граткову релаксацію іонів ²³Na.

Ключові слова: ЯМР, магнітна релаксація, цеоліти, рухливість води, натроліт.

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В Al-натролите (Na₂Al₂Si₃O₁₀ · 2H₂O) и Ga-натролите (Na₂Ga₂Si₃O₁₀ · 2H₂O изучены температурные зависимости спектров ЯМР и MAS ЯМР ядер ²³Na. Было показано, что в Al-натролите и Ga-натролите при T < 400 К отсутствует диффузия ионов натрия. Получены температурные зависимости времени спин-решеточной релаксации времени T_1 в Al-натролите и Ga-натролите. Обсуждается влияние молекулярной подвижности воды в наноканалах структуры натролита на спин-решеточную релаксацию ионов ²³Na.

Ключевые слова: ЯМР, магнитная релаксация, цеолиты, подвижность воды, натролит.

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